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# Controlled crack propagation for atomic precision handling of wafer-scale two-dimensional materials

Jaewoo Shim<sup>1,2\*</sup>, Sang-Hoon Bae<sup>1,2\*</sup>, Wei Kong<sup>1,2\*</sup>, Doyoon Lee<sup>1,2\*</sup>, Kuan Qiao<sup>1,2</sup>, Daniel Nezhich<sup>3</sup>, Yong Ju Park<sup>4</sup>, Ruike Zhao<sup>1,5</sup>, Suresh Sundaram<sup>6</sup>, Xin Li<sup>6</sup>, Hanwool Yeon<sup>1,2</sup>, Chanyeol Choi<sup>1,2</sup>, Hyun Kum<sup>1,2</sup>, Ruoyu Yue<sup>7</sup>, Guanyu Zhou<sup>7</sup>, Yunbo Ou<sup>8</sup>, Kyusang Lee<sup>1,2,9</sup>, Jagadeesh Moodera<sup>8</sup>, Xuanhe Zhao<sup>1</sup>, Jong-Hyun Ahn<sup>4</sup>, Christopher Hinkle<sup>7,10</sup>, Abdallah Ougazzaden<sup>6</sup>, Jeehwan Kim<sup>1,2,11,12†</sup>

<sup>1</sup>Mechanical Engineering, Massachusetts Institute of Technology, Cambridge, MA, USA. <sup>2</sup>Research Laboratory of Electronics, Massachusetts Institute of Technology, Cambridge, MA, USA. <sup>3</sup>Lincoln Laboratory, Massachusetts Institute of Technology, Lexington, MA, USA. <sup>4</sup>School of Electrical and Electronic Engineering, Yonsei University, Seoul 03722, Republic of Korea. <sup>5</sup>Department of Mechanical and Aerospace Engineering, The Ohio State University, Columbus, OH 43210, USA. <sup>6</sup>School of Electrical and Computer Engineering, Georgia Institute of Technology, UMI 2958 GT-CNRS, GT-Lorraine, Metz, France. <sup>7</sup>Materials Science and Engineering, University of Texas at Dallas, Richardson, TX, USA. <sup>8</sup>Department of Physics, and Plasma Science and Fusion Center, Massachusetts Institute of Technology, Cambridge, MA, USA. <sup>9</sup>Departments of Electrical & Computer Engineering and Materials Science Engineering, University of Virginia, Charlottesville, VA, USA. <sup>10</sup>Department of Electrical Engineering, University of Notre Dame, Notre Dame, IN, USA. <sup>11</sup>Materials Science and Engineering, Massachusetts Institute of Technology, Cambridge, MA, USA. <sup>12</sup>Microsystem Technology Laboratory, Massachusetts Institute of Technology, Cambridge, MA, USA.

\*These authors contributed equally to this work.

†Corresponding author. Email: jeehwan@mit.edu

**Although flakes of two-dimensional (2D) heterostructures at micrometer scale can be formed with adhesive-tape methods, isolation of 2D flakes into monolayers is extremely time-consuming as it is a trial-and-error process. Controlling the number of 2D layers through direct growth also presents difficulty because of the high nucleation barrier on 2D materials. We demonstrate a layer-resolved 2D material splitting technique that permits the high-throughput production of multiple monolayers of wafer-scale (5 centimeter diameter) 2D materials by splitting single stacks of thick 2D materials grown on a single wafer. Wafer-scale uniformity of *h*-BN, WS<sub>2</sub>, WSe<sub>2</sub>, MoS<sub>2</sub>, and MoSe<sub>2</sub> monolayers was verified by photoluminescence (PL) response and by substantial retention of electronic conductivity. We fabricated wafer-scale van der Waals heterostructures, including field-effect transistors, with single-atom thickness resolution.**

Because of their in-plane stability and weak out-of-plane interaction, two-dimensional (2D) materials can be stacked together to form a multitude of device types with a broad spectrum of functionalities (1–6). Construction of 2D material-based heterostructures is often described as stacking Lego blocks (4). In order to tailor the 2D heterostructure characteristics for specific functionalities, it is essential to be able to isolate 2D materials into monolayer films and stack them with monolayer precision. The most common method for assembling these 2D blocks is by using the adhesive-tape method (7), where stacking of micrometer-scale flakes has been demonstrated (4). However, this method cannot reliably produce monolayer 2D crystals from bulk materials. The process becomes more complicated if the heterostructure design requires several different types of 2D material monolayers. Multiple monolayer flakes must be initially secured for each 2D material, which becomes extremely time consuming. Moreover, although the isolation of flakes into nominal monolayer has been demonstrated, the lateral dimensions (hundreds of micrometers) are not sufficient to guarantee the fabrication of large-scale 2D heterostructures (8). In parallel,

numerous efforts have been made to directly grow 2D heterostructures at the wafer scale (6, 9, 10). Recently, metal-organic chemical vapor deposition (MOCVD) growths of wafer-scale monolayer 2D materials and their heterostructures have been successfully demonstrated for some of transition metal dichalcogenides (TMDCs) at a specific growth condition (11, 12).

We introduce a layer-resolved splitting (LRS) technique that can be universally applied to produce 2D material monolayers at the wafer scale. This method requires one short growth of thick 2D materials on the wafer at a relaxed growth condition and subsequently harvests the multilayers to individual monolayers through a wafer-scale splitting process. This method allows for the high-throughput production of monolayer 2D materials with single-atom thickness precision for the fabrication of wafer-scale van der Waals (vdW) heterostructures. We demonstrated the wafer-scale LRS for various monolayers of 2D materials, including hexagonal boron nitride (*h*-BN), tungsten disulfide (WS<sub>2</sub>), tungsten diselenide (WSe<sub>2</sub>), molybdenum disulfide (MoS<sub>2</sub>), and molybdenum diselenide (MoSe<sub>2</sub>). These films are readily stackable for

forming wafer-scale 2D heterostructures. We have revealed the underlying mechanics that allows precise control of crack propagation, enabling LRS of the 2D material multilayers into multiple individual monolayers. The wafer-scale monolayer of TMDCs after LRS exhibits substantial photoluminescence (PL) enhancement uniformly across a 2-inch (5-cm) wafer at a photon energy corresponding to its monolayer. Through this digital control of monolayer 2D materials, we demonstrate wafer-scale vdW heterostructures with single-atom thickness resolution. The heterostructure devices prepared by LRS and quasi-dry stacking process exhibited substantial improvement in electrical and optical uniformity across the wafer. This finding will open up new opportunities for 2D material research community as it provides a reliable pathway to form wafer-scale 2D heterostructures with monolayer resolution.

The schematic of the LRS process (Fig. 1A) shows a thick 2D material with an arbitrary number of layers grown on sapphire. Because of the difficulty in controlling the nucleation of the 2D material, new nucleation sites inevitably appear on top of the initial nucleation layer on the wafer before full substrate coverage. Thus, growths of multilayer 2D materials result in irregular, discontinuous films at the top, but does leave uniform continuous films underneath. Once the entire multilayer film is removed from the sapphire wafer, the continuous 2D material films on the bottom can be split into many monolayers with the LRS process. We designed the LRS process based on differences in the interfacial toughness ( $\Gamma$ ) of these materials. We used a 600-nm-thick nickel (Ni) film as an atomic-scale adhesive because the reported  $\Gamma$  between 2D materials and Ni ( $\Gamma_{2D-Ni} \sim 1.4 \text{ J m}^{-2}$ ) (13) is three times greater than that of the vdW interface between layers in 2D materials ( $\Gamma_{2D-2D} = 0.45 \text{ J m}^{-2}$ ) (14). The  $\Gamma$  between 2D materials and sapphire ( $\Gamma_{2D-Sapphire}$ ) has been empirically deduced (see methods) and its average value ( $\Gamma_{2D-Sapphire} = 0.26 \text{ J m}^{-2}$ ) is less than  $\Gamma_{2D-2D}$ . Application of a bending moment during the liftoff of a Ni/2D material stack on sapphire supplies elastic strain energy to the bottom interfaces (Fig. 1B). The strain energy per unit area is released upon delamination when reaching to desired strain release rate ( $G$ ). As shown in Fig. 1C, an external force creates a bending moment, resulting in a torque applied across a small distance in the sample that initiates spalling mode fracture where the cracks propagate downward due to mixed mode I and mode II fracture (15). 2D materials grown at the edge of the wafer is typically defective so that the crack propagation can be facilitated. Spalling mode fracture occurs due to the external bending moment applied. In addition to the opening mode stress (mode I) acting on the crack tip, a shear field (mode II) is created to guide the cracks into the bulk. Thus, exfoliation of Ni/2D material stacks can separate the 2D-sapphire interface that has the weakest interfacial toughness and allow clean separation of

2D materials from the wafer.

We first applied the LRS technique to 4-nm-thick  $\text{WS}_2$  grown on a sapphire wafer through vapor phase epitaxy (VPE), characterizing each step of the process. As shown in the Raman mapping of as-grown  $\text{WS}_2$  on the sapphire wafer, the  $E_{2g}^1$  peak of the  $\text{WS}_2$  indicates that the sapphire wafer is fully covered with  $\text{WS}_2$  after the growth (see Fig. 1D and fig. S1). Atomic force microscopy (AFM) topology taken from the top of the as-grown  $\text{WS}_2$  (Fig. 1E) shows that a rough discontinuous surface morphology developed from uncontrollable nucleation during growth, which is typically observed for the as-grown 2D materials. The LRS process was initiated with an exfoliation step to remove the entire  $\text{WS}_2$  stack from the sapphire wafer. A second exfoliation step isolated them into monolayer films. For the first exfoliation, a 600-nm-thick Ni film was deposited on multilayer  $\text{WS}_2$  followed by application of thermal release tape on the Ni as a handler. Lifting off the tape/Ni stack successfully separated the weakest  $\text{WS}_2$ -sapphire interface resulting in release of the entire  $\text{WS}_2$  film from the substrate; no sign of  $\text{WS}_2$  was detected from Raman mapping and x-ray photoelectron spectroscopy on the sapphire wafer after exfoliation (see Fig. 1F and figs. S2 and S3). Thus, a moment applied by lift-off supplies strain energy to the Ni/ $\text{WS}_2$  stacks sufficient to delaminate the weakest 2D-sapphire interface (16–18). Successful release of the entire  $\text{WS}_2$  film left pristine  $\text{WS}_2$  layers at the bottom. The bottom  $\text{WS}_2$  layer was continuous and smooth with root mean square (RMS) roughness of 0.5 nm (an average value measured from 15 locations, see Fig. 1G) that was the result of complete merging of the nuclei of the initial layers.

To harvest the continuous  $\text{WS}_2$  monolayer, we deposited a Ni layer on the bottom of the  $\text{WS}_2$  film while retaining the top tape/Ni/ $\text{WS}_2$  stack as-exfoliated (see Fig. 1A). Similar to peeling off the Ni/ $\text{WS}_2$  stack from the sapphire substrate, we applied a moment from the top Ni to initiate spalling mode fracture for guiding the cracks downward. As  $\Gamma_{2D-Ni}$  is substantially higher than  $\Gamma_{2D-2D}$ , the cracks arriving near the bottom Ni propagated through weaker  $\text{WS}_2$ - $\text{WS}_2$  interface directly above the bottom Ni layer (see Fig. 1C). Thus, the Ni/ $\text{WS}_2$  stack separated upon peeling while the bottom Ni strongly adhered to the  $\text{WS}_2$  monolayer, which left a monolayer of  $\text{WS}_2$  on the bottom Ni layer.

We transferred this monolayer film onto an 8-inch (20.3-cm) Si wafer coated with 90-nm of  $\text{SiO}_2$  to investigate the thickness and quality of the exfoliated monolayer  $\text{WS}_2$  film (see fig. S4 for photographs of the entire LRS processes). As shown in Fig. 2, A and B, the 2-inch wafer-scale  $\text{WS}_2$  transferred intact onto the 8-inch  $\text{SiO}_2/\text{Si}$  wafer without prominent wrinkle, folding, or ripple. The average thickness of the transferred  $\text{WS}_2$  is 0.7 nm (average value obtained from scans on 10 locations, Fig. 2C). The exact value for the monolayer is 0.616 nm (19), and we speculate that the greater observed

thickness arose tip-sample interactions under non-contact measurement mode. Successful isolation of WS<sub>2</sub> monolayer was confirmed by the substantial enhancement of peak intensity of photoluminescence (PL) spectra (Fig. 2D) at its direct gap of 1.99 eV, versus the weak and wide PL characteristic of a thick WS<sub>2</sub> layer at its indirect gap of 1.97 eV (20–25). Also, the PL intensity was not further degraded (fig. S5), implying that WS<sub>2</sub> was already degraded before the process given the prolonged time period after growth, which is typically observed in TMDCs after air exposure (26). Moreover, the strong PL intensity originating from the monolayer isolation was uniform across the entire 2-inch wafer area (spatial resolution of 2 mm, Fig. 2E). PL mapping with greater spatial resolution further confirmed uniform monolayer thickness obtained by LRS (see Fig. 2F for 50 μm resolution map on a 1 mm × 1 mm area and fig. S6 for 1 μm resolution maps on 20 μm × 20 μm area). Wafer-scale monolayer thickness was also confirmed by mapping the PL peak position where peaks are all concentrated at its direct gap of 1.99 eV (see fig. S7) (20–25).

All LRS WS<sub>2</sub> films were transferred with a quasi-dry process, where the WS<sub>2</sub> split from the multilayers was directly dry-bonded to the SiO<sub>2</sub>-coated Si wafer followed by etching to remove the Ni (16, 17, 27). The average RMS roughness was 0.5 nm after LRS measured from 15 different spots in each material (see fig. S8). We speculate that perfect atomic smoothness (RMS < 0.3 nm) was not obtained because of the polycrystalline nature of the films with a grain size of ~60 nm (see fig. S9). Larger-scale morphological inspection of monolayer 2D materials after LRS with Raman microscopy and laser scanning confocal microscopy revealed no prominent surface defects and macroscopic deformation such as wrinkle, folding, or ripple (see figs. S10 and S11). More importantly, LRS did not substantially degrade electrical quality of the 2D materials. The average Hall mobility measured from few layers before LRS was 106.8 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> versus was 89.5 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> for the monolayer after LRS (see fig. S12) (28).

The LRS process was repeated to harvest additional continuous monolayers until the split layer was no longer continuous. In parallel, an entire 4-nm-thick WS<sub>2</sub> film was transferred from another sapphire substrate onto a 90-nm-thick SiO<sub>2</sub>-coated Si as a reference. Optical micrograph and scanning electron microscope (SEM) image seen in Fig. 3, A and B, showed its non-uniform thickness. Repetition of LRS yielded continuous production of uniform continuous monolayer WS<sub>2</sub> (see fig. S13 for Raman spectra showing right shift for  $E_{2g}^1$  peak compared to that of as-exfoliated 4-nm-thick WS<sub>2</sub> films) through three cycles, as shown in Fig. 3, C to H. After the third cycle, residual WS<sub>2</sub> was observed as discontinuous triangle domains on the top Ni film (see Fig. 3, I and J), indicating that the LRS process has reached the final top layer where nucleated islands failed to merge during growth.

Furthermore, we measured strong direct gap emission peaks (at 1.99 eV) only from the WS<sub>2</sub> obtained through the first three cycles (see Fig. 3K).

We applied our LRS technique to other 2D materials including MoSe<sub>2</sub>, WSe<sub>2</sub>, MoS<sub>2</sub>, and *h*-BN, and confirmed successful monolayer splitting for all of these materials. As shown in fig. S14, 3-nm-thick *h*-BN grown on sapphire was split into three *h*-BN monolayers. In addition, as shown in fig. S15, monolayer isolation of MoSe<sub>2</sub>, WSe<sub>2</sub>, and MoS<sub>2</sub> after LRS was confirmed by PL measurement (see also figs. S16 to S18 for optical microscope images and Raman spectra of MoSe<sub>2</sub>, WSe<sub>2</sub>, and MoS<sub>2</sub> before and after LRS process). We also characterized these monolayers with AFM, confocal Raman microscopy and laser scanning confocal microscopy (figs. S8, S10, and S11).

We then fabricated various 2D heterostructures. First, we fabricated WS<sub>2</sub>/*h*-BN heterostructures with the monolayer WS<sub>2</sub> sandwiched between *h*-BNs to investigate the effect of wafer-scale encapsulation of TMDCs. Reduction of carrier scattering by surface optical phonon has been observed in the flakes of *h*-BN/TMDC/*h*-BN heterostructures (29). Wafer-scale *h*-BN/TMDC/*h*-BN heterostructures were fabricated with LRS together with the quasi-dry transfer process. Monolayer *h*-BN prepared by LRS was transferred twice on a SiO<sub>2</sub>/Si wafer followed by monolayer transfer of WS<sub>2</sub> on top, and capped by transferring *h*-BN twice for encapsulation (see fig. S19 for the schematic description of 2D heterostructure via LRS process). We observed a substantial enhancement of PL intensity uniformly across the wafer (1 mm × 1 mm PL map shown in Fig. 4A and measured across the wafer with same PL results), compared to that of a WS<sub>2</sub> monolayer quasi-dry transferred on SiO<sub>2</sub> (see Fig. 4B).

When we fabricated *h*-BN/WS<sub>2</sub>/*h*-BN heterostructures using a wet-stacking process by scooping 2D materials from a solution (see method for details of the wet-stacking process), we observed a substantially weakened PL response compared to that of heterostructures prepared by the quasi-dry-stacking process (see Fig. 4C), which were attributed to interference coming from poly(methyl methacrylate) (PMMA) residues at the interface (see fig. S20). Representative PL spectra for monolayer WS<sub>2</sub> on SiO<sub>2</sub>, and double-layer *h*-BN/monolayer WS<sub>2</sub>/double-layer *h*-BN prepared by quasi-dry- and wet-stacking are shown in Fig. 4E. Substantial degradation of the PL intensity was observed for the wet-stacked *h*-BN/WS<sub>2</sub> heterostructure even compared to quasi-dry-transferred WS<sub>2</sub> on SiO<sub>2</sub>. A 15-fold enhancement of PL intensity was observed for *h*-BN/WS<sub>2</sub>/*h*-BN quasi-dry-stacks compared to that for WS<sub>2</sub> wet-stacks on SiO<sub>2</sub> (see Fig. 4D).

Our study revealed wafer-scale enhancement in optoelectronic quality of monolayer WS<sub>2</sub> by wafer-scale *h*-BN encapsulation that should permit fabrication of large-scale optoelectronic devices based on van der Waals



heterostructures. Our quasi-dry stacking method is the key to achieving optoelectronic enhancement in TMDCs by *h*-BN encapsulation as the process preserves the quality of the 2D materials. The optical quality of the TMDCs was not severely affected, although the roughness of the surface slightly increases with increasing layer numbers in the heterostructures (see fig. S21) possibly because of the presence of nanometer-scale Ni residues after quasi-dry stacking process (17).

We also fabricated arrays of 2D heterostructure devices at the wafer scale. These heterostructure arrays exhibited uniformly enhanced device performance across the wafer. LRS fabricated MoS<sub>2</sub> and *h*-BN multilayers were quasi-dry-stacked on a SiO<sub>2</sub>/Si wafer to form field effect transistors (FETs). Deleterious trap charges from the substrate were avoided by using *h*-BN (see Fig. 4F for the structure of our 2D heterostructure device) (30–32). Figure 4G shows an optical image of the 10 × 10 arrays of MoS<sub>2</sub> FETs on a 1-cm<sup>2</sup> wafer. Although both FETs with and without *h*-BN exhibited high on-off ratios of >10<sup>7</sup> (see fig. S22), the FETs without *h*-BN shows very large hysteresis in its drain current-gate voltage sweep which is detrimental to its transistor operation. However, substantial suppression of hysteresis has been observed in FETs with *h*-BN (see Fig. 4H). This reduction of hysteresis was obtained uniformly across the arrays in the wafer. As shown in the map of hysteresis of FETs with and without *h*-BN, the FET arrays with *h*-BN showed excellent reduction in hysteresis during a gate voltage sweep applied uniformly across the wafer (see Fig. 4I).

The merit of quasi-dry stacking process was also highlighted by fabricating arrays of WSe<sub>2</sub>/graphene vertical transistors (see fig. S23 for an optical image of the 8 × 8 arrays of WSe<sub>2</sub>/graphene vertical transistors made on a 1 cm × 1 cm wafer). As shown in the map of on/off ratios of vertical transistors (fig. S23), our transistor arrays made by quasi-dry stacking showed excellent device-to-device uniformity with 9.6% variation (see fig. S24 for the statistical distribution). Vertical transistor arrays fabricated by the wet-stacking process had a device-to-device variation of 26%.

We have demonstrated that the LRS technique can precisely isolate monolayers of 2D materials at the wafer scale from multilayers of 2D materials, thus allowing wafer-scale integration of 2D materials. LRS permits high throughput manufacturing of wafer-scale monolayers of various 2D materials from thick 2D material films grown at relaxed growth conditions. We have demonstrated various 2D heterostructure devices at the wafer by applying LRS together with quasi-dry stacking. These heterostructure devices exhibit uniform device performances across the wafer. This high throughput manufacturing of 2D heterostructures will become the stepping stone for the commercialization of 2D material-based devices.

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## SUPPLEMENTARY MATERIALS

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Materials and Methods

Figs. S1 to S24

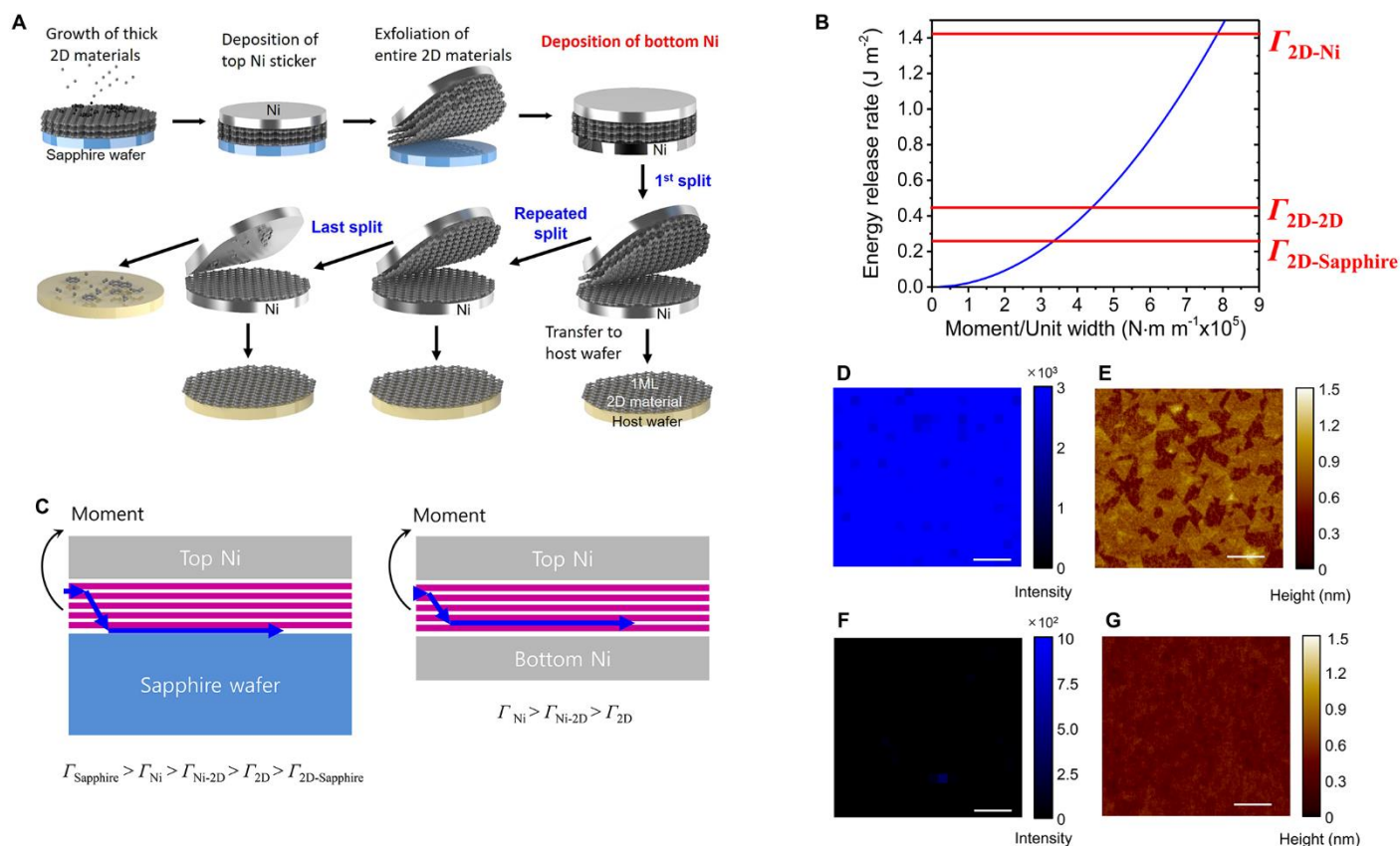
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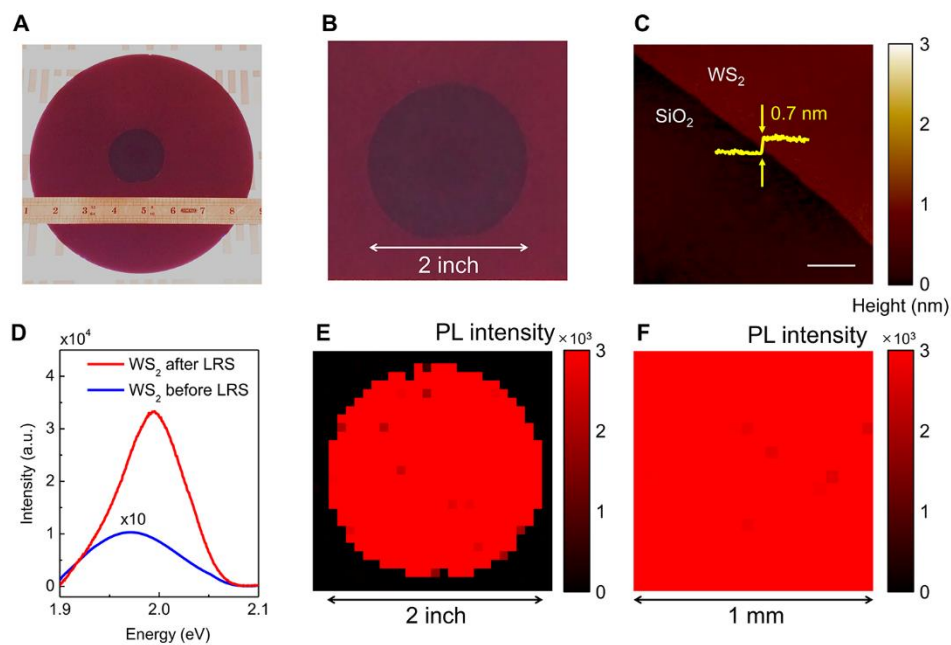
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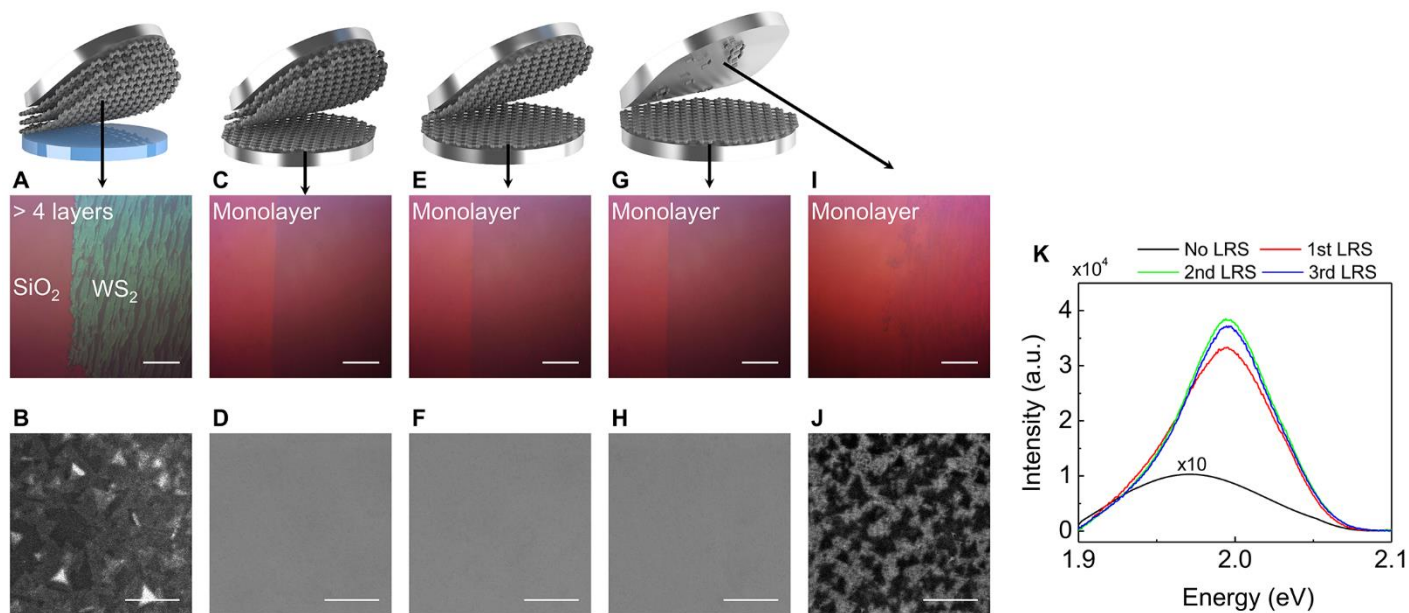


**Fig. 1. Layer-resolved splitting (LRS) of 2D materials.** (A) Schematic illustration explaining the LRS process of 2D materials. (B) modeling of energy release rate according to applied moment. (Ni thickness: 600 nm) (C) Schematics of crack progression during LRS for first exfoliation of entire 2D materials from sapphire wafer (left) and exfoliation of the bottom monolayer 2D material (right). (D) Raman intensity mapping at  $E_{12g}^1$  peak ( $353\text{ cm}^{-1}$ ) of  $\text{WS}_2$  grown on a sapphire substrate, with laser wavelength and power of 532 nm and 2 mW in continuous waveform respectively, where the spatial resolution is  $2\text{ }\mu\text{m}$ . (E) AFM topology taken from the top of as-grown 4-nm-thick  $\text{WS}_2$  on the sapphire wafer. (F) Raman mapping image showing the intensity of  $E_{12g}^1$  peak ( $353\text{ cm}^{-1}$ ) on sapphire substrate after exfoliation of  $\text{WS}_2$  layer, with laser wavelength and power of 532 nm and 2 mW, respectively. (G) AFM topology taken from the bottom of  $\text{WS}_2$  layer after exfoliation. Scale bars for Raman mapping images and AFM topology images are  $2\text{ }\mu\text{m}$  and  $50\text{ nm}$ , respectively.



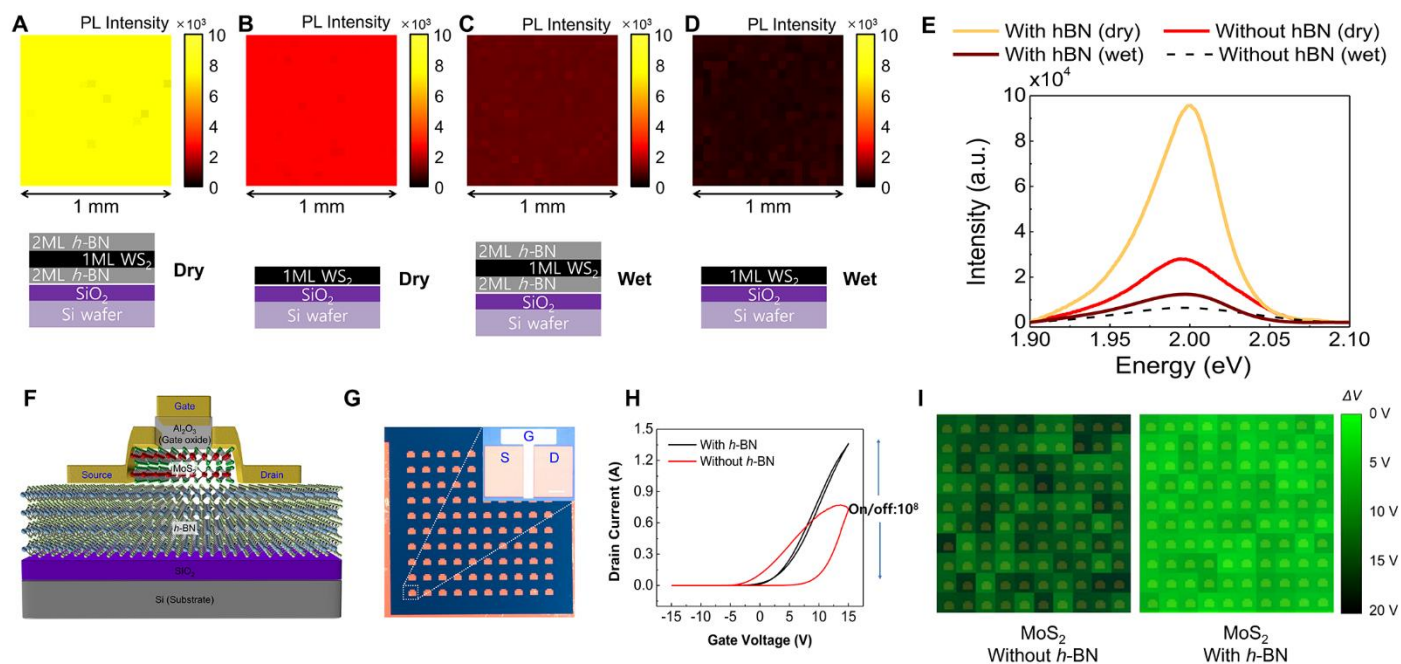
**Fig. 2. Wafer-scale monolayer 2D material obtained by LRS process.** (A) Optical image of 2-inch wafer-scale WS<sub>2</sub> monolayer on 8-inch oxidized Si wafer obtained through the LRS process. (B) Macrograph of 2-inch wafer-scale WS<sub>2</sub> monolayer. (C) AFM image and height profile of the WS<sub>2</sub> monolayer transferred on SiO<sub>2</sub>/Si wafer. Scale bar: 1 μm. (D) Representative PL spectra of as-exfoliated 4-nm-thick WS<sub>2</sub> (blue solid line) and monolayer WS<sub>2</sub> obtained by LRS process (red solid line), where the PL spectrum for monolayer WS<sub>2</sub> is multiplied by 10 times to see the clear peak position. (E) Wafer-scale PL mapping image at 1.99 eV of WS<sub>2</sub> on the SiO<sub>2</sub>/Si wafer, where the spatial resolution is 2 mm. (F) Large-scale (1 mm × 1 mm) PL intensity map from 1.99 eV for WS<sub>2</sub> on SiO<sub>2</sub>, where the spatial resolution is 50 μm (see fig. S6 for 1 μm resolution map). All PL spectra were taken at the same laser power (2 mW) and wavelength (532 nm).





**Fig. 3. Split of thick 2D materials into many monolayers via LRS process and their characterization.** (A, C, E, G, and I) Optical micrographs and (B, D, F, H, and J) plan-view SEM images for as-exfoliated thick WS<sub>2</sub> (A) and (B), monolayers of WS<sub>2</sub> obtained by first (C) and (D), second (E) and (F), third (G) and (H), and the last (I) and (J) LRS processes, respectively. They are all transferred on 90-nm-thick SiO<sub>2</sub>/Si substrates. (K) Representative PL spectra for as-exfoliated thick WS<sub>2</sub> (black solid line) multiplied by ten times to see the clear peak position, monolayers of WS<sub>2</sub> obtained by first (red solid line), second (green solid line), third (blue solid line) LRS processes. All PL spectra were taken at the same laser power (2 mW) and wavelength (532 nm). Scale bars for optical microscopy images and plan-view SEM images are 50 μm and 400 nm, respectively.





**Fig. 4. Wafer-scale 2D heterostructures.** (A) PL mapping images at 1.99 eV of double-layer (2 ML) *h*-BN/monolayer (1 ML) WS<sub>2</sub>/2 ML *h*-BN heterostructure made by quasi-dry-stacking, (B) 1 ML WS<sub>2</sub> on SiO<sub>2</sub> made by quasi-dry-transfer, (C) 2 ML *h*-BN/1 ML WS<sub>2</sub>/2 ML *h*-BN heterostructure made by wet-stacking, and (D) 1 ML WS<sub>2</sub> on SiO<sub>2</sub> made by wet-stacking. (E) Representative PL spectra of all different structures made by different methods. All PL spectra were taken at the same laser power (2 mW) and wavelength (532 nm). (F) Schematic of a MoS<sub>2</sub>-based field effect transistor (FET). (G) Macrograph of 10 × 10 FET arrays integrated on a SiO<sub>2</sub>/Si wafer with 1 cm × 1 cm size. Inset shows micrograph of an individual device. Scale bar: 100 μm. The device area is defined by the gap between source and drain electrodes. (H) Representative drain current–gate voltage ( $I_D$ – $V_G$ ) characteristics of MoS<sub>2</sub>-based FETs at drain voltage ( $V_{DS}$ ) = 1 V. (I) 2D color maps of hysteresis voltage extracted from  $I_D$ – $V_G$  curves at  $V_{DS}$  = 1 V in transistor arrays made without *h*-BN (left) and with *h*-BN (right).

## Controlled crack propagation for atomic precision handling of wafer-scale two-dimensional materials

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